

Bokarev *et al.* Reply: In the Comment [1] Föhlisch *et al.* claim that our recently proposed model explaining the role of electron delocalization into the solvation shell after core-level excitation [2] is not substantiated by the provided experimental and theoretical data. The Comment is based on three aspects: the nature of the core-hole excited state, the interplay between time scales of electron delocalization and radiative or nonradiative core-hole decay, and the role of x-ray optical effects. In the second comment [3], Green *et al.* put forward a simpler explanation of the observed differences between partial fluorescence yield (PFY) and x-ray absorption spectroscopy (XAS) L -edge spectra. They invoke ligand field multiplet theory to conclude that the differences are “consequences of the different formalisms which govern PFY and XAS.”

Föhlisch *et al.* endeavor to invalidate our findings arguing that state-dependent fluorescence yield is an atomic effect without relevant electron delocalization. Their conclusion has been drawn on the basis of spectroscopic and theoretical investigations of $\text{Cr}^{3+}(\text{aq})$ [4]. Reference [4] does, however, lack a discussion of the underlying resonant x-ray emission spectra, which allow for tracking orbital mixing by experiment as we have recently demonstrated for $\text{Fe}^{2+}(\text{aq})$ [5]. Furthermore, recent resonant photoelectron studies of $\text{Fe}^{3+}(\text{aq})$ [6] and $\text{Co}^{2+}(\text{aq})$ [7] have revealed orbital mixing between metal ion and surrounding water molecules, associated with electron delocalization upon core-hole excitation. The effect is especially pronounced for $\text{Fe}^{3+}(\text{aq})$ in which case strong intensity enhancement at the energetic position of the water $3a_1$ orbital occurs [6].

In their multiconfigurational second-order perturbation theory restricted active space (RASPT2) and restricted active space self-consistent field (RASSCF) calculations, Wernet *et al.* [4] included the influence of the ligands beyond the crystal field electrostatic effect, i.e., one-electron orbital mixing. Their RASPT2 calculations account for ligand-metal electron correlation via excitations from and to ligand orbitals. The relevance of these correlation and exchange effects between metal and ligands is a clear manifestation of the nonatomic character of the wave function. Nevertheless, even though Wernet *et al.* admit the notable orbital mixing and substantial changes upon inclusion of metal-ligand electron correlation (Fig. 1 in [8]), they insist that orbital delocalization does not play a role in interpreting the experimental data. The arguments of [3,4] are based on atomic ligand multiplet theory [9]. It is interesting that our data can be reproduced by a well-established semiempirical theory, which is atomic in nature and contains a number of empirical parameters that cannot be consistently obtained from *ab initio* theory. However, the fact that the data can be fitted by an equation that is *a priori* biased towards an atomic model cannot be taken as a falsification of our unbiased *ab initio* methodology.

Although the methodology and results in [4] are correct, the interpretation leaves room for discussion as it is based

on the isolated case of Cr^{3+} , and therefore cannot be considered as general. Wernet *et al.* do not provide the careful analysis of states that we present for the first time. Specifically, we suggest a system-dependent interpretation showing that state-dependent effects on fluorescence decrease in a series $\text{Fe}^{3+} > \text{Co}^{2+} > \text{Fe}^{2+}$ [2], and local atomic and intermolecular effects could be more intertwined [5,10]. As shown in [2], the RASSCF wave function has a complex multiconfigurational nature. To allow for simple and straightforward interpretation, we used a reduced representation of the nature of the states by means of orbital occupation numbers. Occupancies presented in Figs. 1 and S2 in [2] show the difference in core-excited and ground state occupation numbers, thus illustrating the localization of the core-excited electron. We emphasized the separation between the fraction of $2p \rightarrow 3d(t_{2g})$ transitions (prepeak of L_3) and the rest of the L_3 band, which has a more complex nature [1]. This allows for looking specifically at t_{2g} localized states separated from more intense $2p \rightarrow e_g$ transitions. In this respect $\text{Fe}^{3+}(\text{aq})$ can be clearly distinguished from $\text{Fe}^{2+}(\text{aq})$ and $\text{Co}^{2+}(\text{aq})$. We suggest an interpretation where the isolated t_{2g} states in $\text{Fe}^{3+}(\text{aq})$ exhibit different behavior when compared to the rest of the L_3 edge due to stronger electron delocalization.

In previous works [4,8,11] delocalization has been rationalized within a stationary picture where electronic relaxation effects (like orbital mixing) are completed. In [2], we instead shift the focus on the time-dependent interpretation, following recent works by Cederbaum *et al.* [12,13], who showed that charge migration after the instantaneous ionization or interatomic Coulombic decay takes place on a time scale comparable to the core-excited state lifetime of a few femtoseconds. This suggests the importance of the interplay between delocalization and radiative or nonradiative core-excited state decay, which is governed by the ratio of the respective time scales.

As correctly pointed by Föhlisch *et al.* [1], the total lifetime of the core-excited state depends mostly on the nonradiative (Auger-type) decay rate, whereas radiative decay is a minor channel. However, the particular decay channel is not essential for our interpretation. To quantify our argument we have fitted the total lifetime broadening of RASSCF results to the partial-electron-yield x-ray absorption spectrum of $\text{Fe}^{3+}(\text{aq})$ [6]. We find that the lifetime of isolated t_{2g} states in the 707–709 eV spectral region exceeds that of the e_g states above 709 eV by 3 fs. As a possible reason for the longer lifetime, we consider the variation of Auger decay rate due to the state-dependent $2p3d3d$ channel [14]. Based on this lifetime estimate and the characteristic time scales for electron wave packet dynamics, we conclude that the delocalization will be more pronounced for t_{2g} than for e_g states.

Finally, Föhlisch *et al.* [1] anticipate strong polarization effects similar to the solid state phase [15]. However, these findings cannot be readily transferred to the present

solution-phase situation [10,16–19] (see discussion in Supplemental Material [20]).

In summary, we do not agree with the criticism put forward in Refs. [1,3]. Based on our multispectroscopic experimental approach and a comprehensive *ab initio* theory, we demonstrated that state-dependent electron delocalization is a possible electron pathway of the $2p$ core-excited state at the L_3 edge of aqueous transition metal ions.

S. I. Bokarev,¹ M. Dantz,^{2,3} E. Suljoti,^{2,3} K. Atak,^{2,3}
B. Winter,² O. Kühn¹ and E. F. Aziz^{2,3}

¹Institut für Physik

Universität Rostock

Universitätsplatz 3, D-18055 Rostock, Germany

²Joint Ultrafast Dynamics Lab in Solutions and at Interfaces (JULiq)

Helmholtz-Zentrum Berlin für Materialien und Energie

Albert-Einstein-Strasse 15, D-12489 Berlin, Germany

³Freie Universität Berlin

FB Physik, Arnimallee 14, D-14195 Berlin, Germany

Received 6 December 2013; published 28 March 2014

DOI: 10.1103/PhysRevLett.112.129303

PACS numbers: 31.15.vj, 31.70.Dk, 32.80.Aa, 78.70.En

- [1] A. Föhlisch, F. M. F. de Groot, M. Odellius, S. Techert, and P. Wernet, this issue, *Phys. Rev. Lett.* **112**, 129302 (2014).
- [2] S. I. Bokarev, M. Dantz, E. Suljoti, O. Kühn, and E. F. Aziz, *Phys. Rev. Lett.* **111**, 083002 (2013).
- [3] R. J. Green, D. Peak, A. J. Achkar, J. S. Tse, A. Moewes, D. G. Hawthorn, and T. Z. Regier, preceding Comment, *Phys. Rev. Lett.* **112**, 129301 (2014).
- [4] P. Wernet, K. Kunnus, S. Schreck, W. Quevedo, R. Kurian, S. Techert, F. M. F. de Groot, M. Odellius, and A. Föhlisch, *J. Phys. Chem. Lett.* **3**, 3448 (2012).
- [5] K. Atak, S. I. Bokarev, M. Gotz, R. Golnak, K. M. Lange, N. Engel, M. Dantz, E. Suljoti, O. Kühn, and E. F. Aziz, *J. Phys. Chem. B* **117**, 12613 (2013).
- [6] S. Thürmer, R. Seidel, W. Eberhardt, S. E. Bradforth, and B. Winter, *J. Am. Chem. Soc.* **133**, 12528 (2011).
- [7] R. Seidel, S. Ghadimi, K. M. Lange, S. Bonhommeau, M. A. Soldatov, R. Golnak, A. Kothe, R. Könnecke, A. Soldatov, S. Thürmer, B. Winter, and E. F. Aziz, *J. Am. Chem. Soc.* **134**, 1600 (2012).
- [8] I. Josefsson, K. Kunnus, S. Schreck, A. Föhlisch, F. de Groot, P. Wernet, and M. Odellius, *J. Phys. Chem. Lett.* **3**, 3565 (2012).
- [9] F. de Groot and A. Kotani, *Core Level Spectroscopy of Solids* (CRC Press, Boca Raton, 2008).
- [10] E. Suljoti, R. Garcia-Diez, S. I. Bokarev, K. M. Lange, R. Schoch, B. Dierker, M. Dantz, K. Yamamoto, N. Engel, K. Atak, O. Kühn, M. Bauer, J.-E. Rubensson, and E. F. Aziz, *Angew. Chem., Int. Ed.* **52**, 9841 (2013).
- [11] L.-Å. Näslund, M. Cavalleri, H. Ogasawara, A. Nilsson, L. G. M. Pettersson, P. Wernet, D. C. Edwards, M. Sandström, and S. Myneni, *J. Phys. Chem. A* **107**, 6869 (2003).
- [12] A. I. Kuleff, and L. S. Cederbaum, *Phys. Rev. Lett.* **106**, 053001 (2011).
- [13] A. I. Kuleff, and L. S. Cederbaum, *Phys. Rev. Lett.* **98**, 083201 (2007).
- [14] F. M. F. de Groot, M. A. Arrio, P. Sainctavit, C. Cartier, and C. T. Chen, *Solid State Commun.* **92**, 991 (1994).
- [15] R. Kurian, K. Kunnus, P. Wernet, S. M. Butorin, G. Piete, and F. M. F. de Groot, *J. Phys. Condens. Matter* **24**, 452201 (2012).
- [16] M. Odellius, H. Ogasawara, D. Nordlund, O. Fuchs, L. Weinhardt, F. Maier, E. Umbach, C. Heske, Y. Zubavichus, M. Grunze, J. D. Denlinger, L. G. M. Pettersson, and A. Nilsson, *Phys. Rev. Lett.* **94**, 227401 (2005).
- [17] J. Forsberg, J. Gråsjö, B. Brena, J. Nordgren, L.-C. Duda, and J.-E. Rubensson, *Phys. Rev. B* **79**, 132203 (2009).
- [18] T. Tokushima, Y. Horikawa, H. Arai, Y. Harada, O. Takahashi, L. G. M. Pettersson, A. Nilsson, and S. Shin, *J. Chem. Phys.* **136**, 044517 (2012).
- [19] B. Dierker, E. Suljoti, K. Atak, K. M. Lange, N. Engel, R. Golnak, M. Dantz, K. Hodeck, M. Khan, N. Kosugi, and E. F. Aziz, *New J. Phys.* **15**, 093025 (2013).
- [20] See Supplemental Material at <http://link.aps.org/supplemental/10.1103/PhysRevLett.112.129303> for discussion of polarization effects on XAS and RIXS spectra of transition metal compounds in solution.